

# Radial Probability Density

$$\int dV |\Psi|^2 = 1$$

$$\text{Prob (find part. in } dV \text{ about } \vec{r}) = |\Psi(\vec{r})|^2 dV$$

$$\text{If } l=0, \Psi = \Psi(r) \text{ then } \int dV |\Psi|^2 = \int dr 4\pi r^2 |\Psi(r)|^2$$

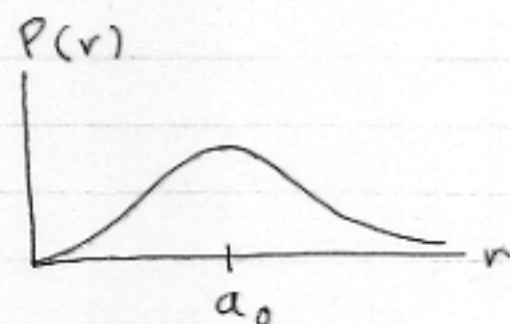
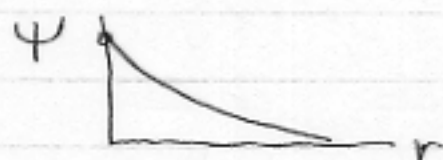
$$\text{Prob (find in } r \rightarrow r+dr) = \boxed{P(r) dr = 4\pi r^2 |\Psi(r)|^2 dr}$$

$P(r)$  = radial probability density

$$\text{Ground state: } \Psi_{100} = A e^{-r/a_0}$$

$$P(r) = |A|^2 4\pi r^2 e^{-2r/a_0}$$

Notice  $P(r)$  very different from  $\Psi(r)$ :



If  $l \neq 0$ ,  $\Psi = \Psi(r, \theta, \phi) = R(r) Y(\theta, \phi)$ , then

$$\int dV |\Psi|^2 = \underbrace{\int dr r^2 |R|^2}_1 \underbrace{\int d\Omega |Y|^2}_{\text{"solid angle"} = 1} = 1$$

$$\text{Prob (find in } r \rightarrow r+dr) = r^2 |R|^2 dr$$

$$\boxed{P(r) = r^2 |R|^2} \text{ even if } l \neq 0$$

$$\text{Note: if } \Psi = \Psi(r) = R \cdot Y = R \cdot \frac{1}{\sqrt{4\pi}} \Rightarrow |R|^2 = 4\pi |\Psi|^2$$

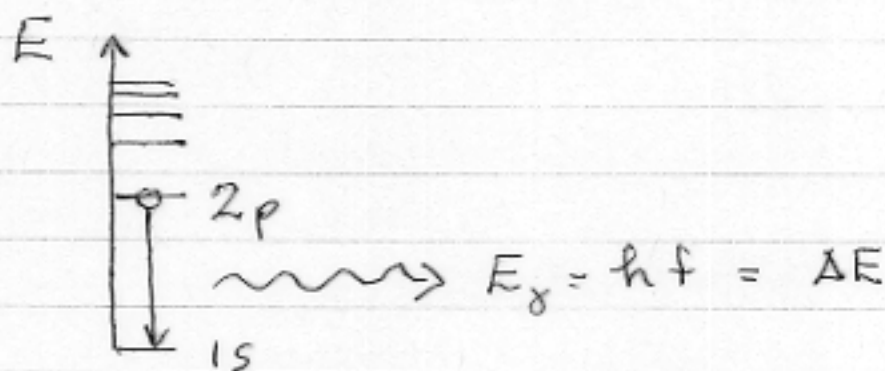
$$\text{so } P(r) = r^2 |R|^2 = 4\pi r^2 |\Psi|^2$$

## H-atom and emission/absorption of radiation:

If H-atom is in excited state ( $n=2, l=1, m=0$ ) then it is in energy eigenstate = stationary state. If atom is isolated, then atom should remain in state  $\Psi_{210}$  forever, since stationary state has simple time dependence:

$$\Psi(\vec{r}, t) = \Psi_{210}(\vec{r}) \cdot e^{-iE_2 t/\hbar}$$

But, experimentally, we find that H-atom emits photon and de-excites:  $\Psi_{210} \rightarrow \Psi_{100}$  in  $\approx 10^{-7} \text{ s}$   
 $\rightarrow 10^{-9} \text{ s}$



The reason that the atom does not remain in stationary state is that it is not truly isolated. The atom feels a fluctuating EM field due to "vacuum fluctuations". Quantum Electrodynamics is a relativistic theory of the EM interaction of matter and light. It predicts that the "vacuum" is not "empty" or "nothing" as previously supposed, but is instead a seething foam of virtual photons and other particles. These vacuum fluctuations interact w/ the electron in the H-atom and slightly alter the potential  $V(r)$ . So eigenstates of the coulomb potential are not eigenstate of the actual potential:

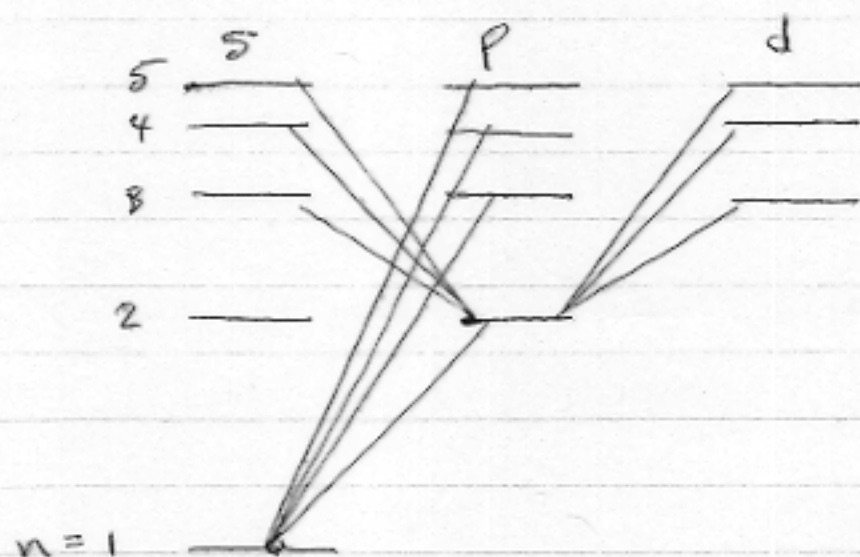
$$V_{\text{coulomb}} + V_{\text{vacuum}}$$

Photons possess an intrinsic angular momentum (spin) of  $1\hbar$ , meaning  $l=1 \Rightarrow |\vec{L}| = \hbar\sqrt{l(l+1)} = \sqrt{2}\hbar$  and  $L_{z, \text{Max}} = \hbar$

So when an atom absorbs or emits a single photon, its angular momentum must change by  $1\hbar$ , by Conservation of Angular Momentum, so the orbital angular momentum quantum nbr  $l$  must change by 1.

"Selection Rule":  $\Delta l = \pm 1$  in any process involving emission or absorption of 1 photon

$\Rightarrow$  allowed transitions are:



If an  $\text{H}^-$  atom is in state  $2s$  ( $n=2, l=0$ ) then it cannot de-excite to ground state by emission of a photon. (since this would violate the selection rule) It can only lose its energy (de-excite) by collision with another atom or via a rare 2-photon process